

QCD saturation and soft processes

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Abstract. We show that an approximate solution to the amended non-linear Balitsky-Kovchegov evolution equation which was formulated for hard (large Q^2) QCD processes, can be extended to provide a good description of photoproduction and soft hadronic (non perturbative) reactions.

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1 Introduction

Further proof that the linear QCD evolution equations fail to describe the DIS data for low values of Q^2 has recently been provided by the ZEUS collaboration [1]. See Fig. 1. For values of $Q^2 \geq 1 \text{ GeV}^2$ the ZEUS NLO QCD fit provides an excellent description of the data, however, as Q^2 becomes smaller the discrepancy between the predictions and the data increases.

The reason for this is well known, and is due to the fact that in the linear evolution equations DGLAP and BFKL, the splitting functions only incorporate the production of partons (gluons). GLR [2] in their classical paper pointed out that when the density of partons becomes sufficiently

large (that they overlap), one has to include non-linear (annihilation) processes in the evolution equations. A recent application of the ideas of GLR using the dipole formalism has been suggested by Balitsky [3] and Kovchegov [4] (which we will denote by BK). The advantage of the BK equation is that:

- (i) it accounts for saturation effects due to high parton densities;
- (ii) it sums higher twist contributions;
- (iii) it allows one to extrapolate to small values of Q^2 (large distances).

2 Amended non-linear BK equation

The BK equation can be written in the form [5]

$$\frac{d\tilde{N}(\mathbf{x}_{01}, b, Y)}{dY} = \frac{C_F \alpha_s}{\pi^2} \int_{\rho} d^2\mathbf{x}_2 \frac{\mathbf{x}_{01}^2}{\mathbf{x}_{02}^2 \mathbf{x}_{12}^2} \cdot (2\tilde{N}(\mathbf{x}_{02}, b, Y) - \tilde{N}(\mathbf{x}_{02}, b, Y) \tilde{N}(\mathbf{x}_{12}, b, Y)) \quad (1)$$

where $\tilde{N}(\mathbf{r}_{\perp}, b, Y)$ denotes the imaginary part of the amplitude of a dipole of size r_{\perp} elastically scattered at impact parameter b , i.e.

$$\tilde{N}(\mathbf{r}_{\perp}, b, Y) = \text{Im} a_{dipole}^{el}(\mathbf{r}_{\perp}, b, Y). \quad (2)$$

$Y = \ln(1/x_{BJ})$ denotes the rapidity. In (1) the linear term corresponds to the LO BFKL evolution, while the non-linear negative term is responsible for unitarization.

A deficiency of the BK equation is that it does not contain evolution in Q^2 , and therefore lacks the correct short distance behaviour. To remedy this we have introduced a correcting function $\Delta\tilde{N}(\mathbf{r}_{\perp}, b, Y)$ which accounts for the DGLAP evolution in Q^2 . Our full solution therefore consists of the sum of two terms

$$N(\mathbf{r}_{\perp}, b, Y) = \tilde{N}(\mathbf{r}_{\perp}, b, Y) + \Delta\tilde{N}(\mathbf{r}_{\perp}, b, Y) \quad (3)$$

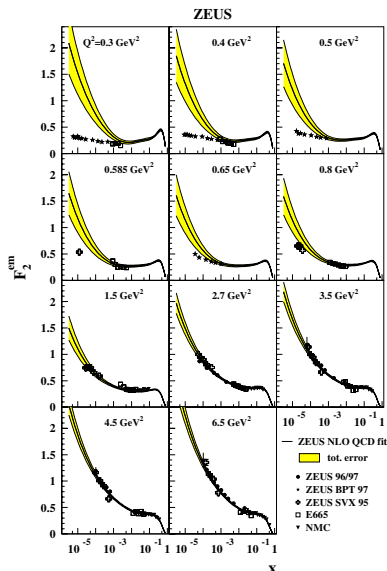


Fig. 1. F_2 data at very low Q^2 compared to ZEUS-S NLO QCD fit, from [1]

Full details of the calculation are contained in [5], where we have simplified the treatment of the dependence on the impact parameter b . We solve for the initial condition at $b = 0$, and then restore b dependence by assuming factorization and that the dipole profile function inside the target is given by a Gaussian distribution $S(b) = e^{-\frac{b^2}{R^2}}$ for all b . This is obviously an over simplification (see discussion in Sect. 5), and the true impact parameter dependence of $N(\mathbf{r}_\perp, b, Y)$ is far more complex.

The deep inelastic structure function F_2 is related to the dipole cross section (defined in (2))

$$F_2(x_{BJ}, Q^2) = \frac{Q^2}{4\pi^2} \int d^2r_\perp \int dz P^{\gamma^*}(Q^2; r_\perp, z) \cdot \sigma_{dipole}(r_\perp, x_{BJ}) \quad (4)$$

Where $P^{\gamma^*}(Q^2; r_\perp, z)$ denotes the probability of the decay of a virtual photon having four momentum Q^2 into a colourless dipole ($q\bar{q}$ pair) of size r_\perp , with the quark (anti-quark) taking a fraction z ($1-z$) of the virtual photon's momentum.

3 Numerical solution of the equation for DIS

In making our fit we include all available data satisfying the following criteria $10^{-7} \leq x_{BJ} \leq x_{0BJ} = 10^{-2}$ and $Q^2 \geq 0.04 \text{ GeV}^2$. We solve (1) as an evolution equation in rapidity with a fixed grid in r_\perp ($= 2/Q \text{ GeV}$) space and a dynamical step in rapidity. We fit to 345 data points for $F_2(x_{BJ}, Q^2)$, and obtain an overall $\chi^2/df \approx 1$. In [5] we also compare to data on the logarithmic slopes of F_2 i.e. $\frac{dF_2}{d\ln Q^2}$ and $\frac{d\ln F_2}{d(\ln 1/x)}$. The values that we obtained for $\lambda = \frac{d\ln F_2}{d(\ln 1/x)}$ were $\lambda \approx 0.07$ for $Q^2 = 0.05 \text{ GeV}^2$ and $x_{BJ} = 10^{-5}$, increasing to $\lambda \approx 0.3$ for $Q^2 = 150 \text{ GeV}^2$ and $x_{BJ} = 4 \cdot 10^{-2}$. i.e. our formalism was successful in describing not only the short range (large Q^2) data but also the ‘‘soft’’ data (at very low Q^2) where the traditional Pomeron intercept is ≈ 0.08 .

4 Extension of BK formalism to photoproduction

The surprising results discussed above, that with the amended BK equation we found an excellent description of all DIS data for $0.05 \leq Q^2 \leq 200 \text{ GeV}^2$ prompted the question, whether this formalism could also successfully describe soft processes e.g. photoproduction?

To extend our formalism to photoproduction [6], it is necessary to make the following alterations:

(i) We need to introduce a finite mass as a cutoff for the r_\perp integration in (4), we take this parameter as m_q (the quark mass).

(ii) The variable x ($= x_{BJ}$) is not defined for γ -p scattering, and we relate x to the energy by introducing a non-perturbative scale Q_0^2 , and taking $x = \frac{Q^2 + Q_0^2}{W_{\gamma p}^2}$. To reduce the number of free parameters we have set $Q_0^2 = 4 m_q^2$.

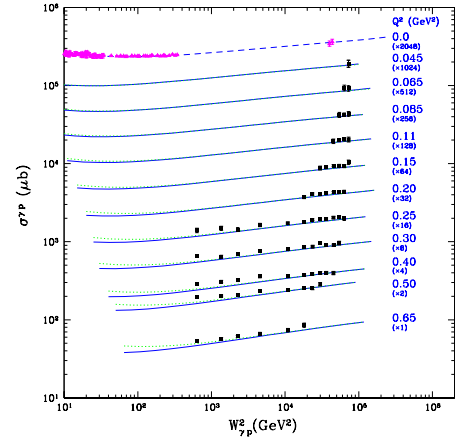


Fig. 2. γ -p and DIS cross-sections at very low Q^2 . *Solid line* VQ and *dotted line* SR parametrization for low energy contribution (see text for details)

On fitting to the high energy photoproduction data we find a value of $m_q = 0.15 \text{ GeV}$. In the colour dipole formalism one can only hope to reproduce the asymptotic energy dependence i.e. the Pomeron contribution. We also need to include a non-singlet (NS) term to account for lower energy (higher x) contributions. We test two possible forms for the NS term:

1) Based on Valence Quark Model (VQ):

$$F_2^{VQ} = \left(\frac{Q^2}{1 \text{ GeV}^2}\right)^{(1+\beta)} \cdot \frac{(1 \text{ GeV}^2 + \mu^2)}{(Q^2 + \mu^2)} \cdot \sum_{i=u,d} e_i^2 q_i^V(Q^2 = 1 \text{ GeV}^2) \quad (5)$$

We freeze the CTEQ6 valence quark contributions at $Q^2 = 1 \text{ GeV}^2$, and on fitting to data, the best fit parameters are $\beta = 0$ and $\mu^2 = 0.13 \text{ GeV}^2$.

2) Based on exchange of Secondary Reggeons (SR):

$$\sigma_{SR}^{\gamma p} = f(0) \frac{\tilde{M}^2}{(Q^2 + \tilde{M}^2)} \cdot \left(\frac{s}{s_0}\right)^{\alpha_R}$$

with $\alpha_R = -0.45$ and $s_0 = 1 \text{ GeV}^2$, $f(0)$ denotes the residue at $Q^2 = 0$. Fit results in $f(0) = 0.19 \text{ mb}$ and $\tilde{M}^2 = 2 \text{ GeV}^2$.

The results of the fit to photoproduction and low Q^2 DIS data with the two alternate low energy contributions are shown in Fig. 2.

5 The BK formalism and soft hadronic processes

Based on the successful extension of our BK formalism to photoproduction (Sect. 4), we also applied the procedure to soft hadronic interactions [7]. This necessitated adapting the basic formula (4) of the dipole model, and hypoth-

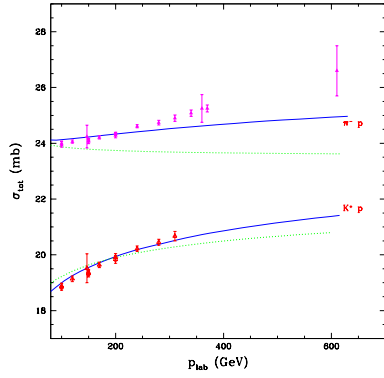


Fig. 3. K^+p and π^-p cross-sections. The *full line* is the prediction in our model and the *dotted line* using the Golec-Biernat Wüsthoff dipole parametrization [10]

esizing that the hadron-proton cross-section is given by:

$$\sigma_{hadron-proton}(x) = \int d^2r_{\perp} |\psi_h(r_{\perp})|^2 \sigma_{dipole}(r_{\perp}, x) \quad (6)$$

where $\psi_h(r_{\perp})$ represents the wave function of the hadron which scatters off the target proton, and the energy dependence is given by $x = \frac{Q_0^2}{W^2}$, with Q_0^2 being an additional non-perturbative scale.

For the hadronic wave functions we use the form suggested by the Heidelberg group, Dosch et al. [8]. The hadronic transverse wave function is taken as a simple Gaussian, where the square of the wave function is given by

$$|\psi_h(r_{\perp})|^2 = \frac{1}{\pi S_M^2} \exp\left(-\frac{r_{\perp}^2}{S_M^2}\right) \quad (7)$$

and S_M is a parameter related to the meson size. We have used $S_{\pi} = 1.08$ fm and $S_K = 0.95$ fm, which is related to the electromagnetic radii. The exotic channel K^+p has no secondary Regge contributions, while for the π^-p reaction we have added a secondary contribution á la Donnachie and Landshoff [9]. Our results are displayed in Fig. 3, for more details regarding the fit shown in Fig. 3 we refer the reader to [7].

For a baryon projectile, we assume that the baryon is constituted of two colour dipoles, one dipole formed around two quarks, and the second dipole from the centre of mass of these two quarks to the third quark in the baryon. Our results for $\bar{p}p$ and pp scattering are shown in Fig. 4. The assumed Gaussian dependence in b , corresponds to a e^{-t} behaviour of the differential cross-section. The full lines in Fig. 4 show that this is a poor description, a much better fit is obtained by assuming that the momentum transfer dependence is of a “form factor” dipole type, which transforms to a Bessel function K_1 dependence in impact parameter space. See [7] for more details.

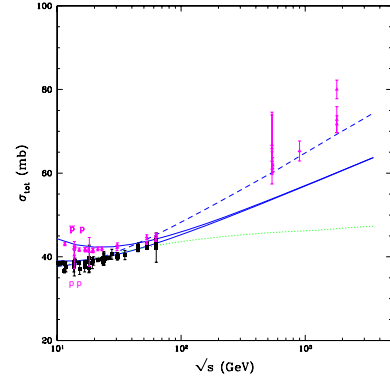


Fig. 4. $\bar{p}p$ and pp total cross-sections. The *full lines* are the model predictions with a Gaussian profile for the impact parameter dependence. The *dashed line* for a K_1 profile and *dotted line* for the GBW dipole model

6 Conclusions

We have shown that our approximate solution to the amended non-linear BK equation which was formulated for hard (large Q^2) processes [5], can be successfully extended to describe photoproduction and soft (non perturbative) hadronic reactions. Two outstanding problems remain:

(i) the impact parameter dependence of the colour dipole. The assumption of factorization in b space, plus imposing a Gaussian like behaviour in b (for all values of b), is obviously naive, as can be judged from the results we have obtained for $\bar{p}p$ scattering. The search for the correct impact parameter dependence of the solution to the BK equations continues.

(ii) the form of the hadronic QCD wave function. The form suggested by [8], is only a first approximation and should be improved.

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